

Crystal and Molecular Structure of 1-Bromotriptycene $\text{BrC}_{20}\text{H}_{13}$

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(Received 28 July 1967)

The crystal structure of 1-bromotriptycene, $\text{BrC}(\text{C}_6\text{H}_4)_3\text{CH}$, has been determined by X-ray diffraction. The space group is $R\bar{3}$. The hexagonal unit cell contains six molecules and has the dimensions $a=11.852 \text{ \AA}$, $c=17.574 \text{ \AA}$. The final value for R is 12%. Lack of better agreement is believed to be due to molecular disorder in the crystal. The bridging C-C bonds have lengths 1.53 \AA at the bromine end and 1.51 \AA at the hydrogen end. These distances and the bond angles at the bridges indicate that there is no conjugation between benzene rings.

Electric dipole moments of isomeric triptycene derivatives (Ogura, Toshiyasu, Kimura, Fujishiro & Nakagawa, 1964) and the infrared spectrum of triptycene (Bartlett & Greene, 1954; Theilacker, Albrecht & Uffmann, 1965) support the conclusion that triptycene is a highly symmetrical and rigid molecule. The ultraviolet spectra of triptycene and a number of its derivatives led Bartlett & Lewis (1950) to conclude that inter-ring resonance occurs. On the other hand, Wilcox (Wilcox, 1960; Wilcox & Craig, 1961) concluded on the basis of quantum mechanical calculations that interaction between benzene rings is negligible. The present investigation was undertaken to determine the length of the bridge bonds and thereby evaluate their bond character. In addition, 1-bromotriptycene, shown in Fig. 1, cannot form without considerable deviation from normal values of one or more of the angles made by the bridge bonds. This study shows which bonds are strained and by what amount.

Experimental

The 1-bromotriptycene used in this investigation was obtained from Prof. A. Streitwieser, who had received it from Prof. L. Friedman. It was synthesized by the method of Friedman & Logullo (1963). Dr E. Ziegler checked the purity of this material by passing it through a gas-liquid chromatographic column. He found only one peak.

Crystals of 1-bromotriptycene suitable for X-ray analysis were prepared by sublimation at approximately 250° . The crystals were rhombohedral in shape. In polarized light the extinction is symmetrical and the birefringence is high. c -Axis views showed a uniaxial interference figure of positive character.

Refractive indices obtained by immersion methods with white light are only approximate because the crystals dissolved in the immersion liquids. A crystal

lying on a rhombohedral face shows ω and ϵ' ; an edge view shows ω and ϵ (Fig. 2). Approximate values for these indices are: $\omega=1.71$, $\epsilon'=1.74$ and $\epsilon=1.81$.

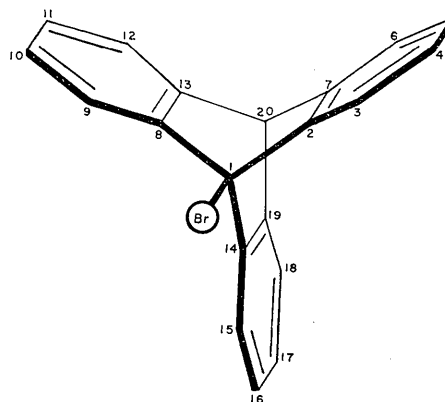


Fig. 1. 1-Bromotriptycene molecule with numbering system used in this investigation. The asymmetric unit is $(\text{BrC}_2)_{1/3}\text{C}_6$.

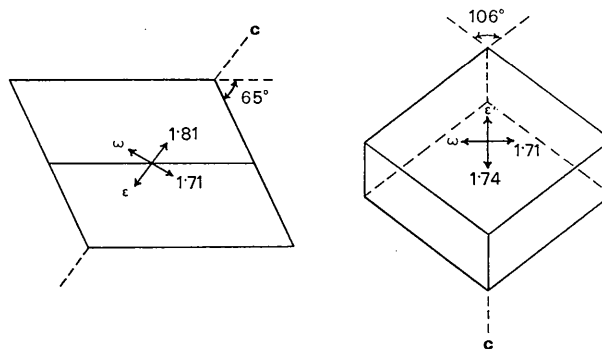


Fig. 2. The directions and magnitudes of the refractive indices ω , ϵ , and ϵ' are shown with respect to the crystal morphology.

Multiple level Weissenberg and precession photographs were taken. The precession photographs are compatible with a Laue symmetry of $\bar{3}$. Reflections were indexed on a hexagonal lattice and were found to obey the rule $(-h+k+l)=3n$. The primitive lattice is therefore rhombohedral and the orientation of the rhombohedral lattice in the hexagonal lattice is obverse.

The observed density of 1.52 g.cm^{-3} requires six molecules per hexagonal cell or two per rhombohedral cell. The most probable space group, therefore, is $R\bar{3}$, and the asymmetric unit is $\text{C}_6\text{H}_4(\text{C}_2\text{HBr})_{1/3}$.

Unit-cell parameters and intensity data were obtained with a diffractometer equipped with a quarter circle goniostat, solid state radiation analyzer and scintillation counter.

2θ scans with a 1° take-off angle for all observable orders of $h0.0$ and $00.l$ led to the following hexagonal unit cell data: $a=11.852(5) \text{ \AA}$, $c=17.574(9) \text{ \AA}$, $V=2,138.0 \text{ \AA}^3$, $Z=6$, ρ calc. $=1.54 \text{ g.cm}^{-3}$, ρ obs. $=1.52 \text{ g.cm}^{-3}$. Estimated standard deviations ($\times 10^3$) are given in parenthesis.

Intensity data were collected by counting for 10 sec with a stationary crystal and stationary counter at a take-off angle of 4° . The background was measured by counting for 10 sec at the minimum between the $K\beta$ and $K\alpha$ peaks. A total of 1050 independent reflections were measured (maximum $2\theta=150^\circ$), 35 of which were assigned zero intensity. Cu $K\alpha$ X-rays ($\lambda=1.5418 \text{ \AA}$) filtered with Ni foil were used. Lorentz and polarization corrections were applied by means of a program developed by A. Zalkin.

The crystal used for intensity measurements had dimensions of $0.088 \text{ mm} \times 0.088 \text{ mm} \times 0.118 \text{ mm}$. The absorption factor for this crystal is 35.0 cm^{-1} and μR is about 0.21 or less, making the absorption quite low. No corrections were applied to the data for either absorption or extinction.

Determination of the structure

The three-dimensional Patterson function showed many well-resolved peaks. Since the 1-bromotriptycene molecule has threefold symmetry and there are six molecules in the hexagonal cell, the bromine and the two aliphatic carbon atoms were assumed to occupy the special positions (c) with parameters $(0, 0, \pm z)$, $(\frac{2}{3}, \frac{1}{3}, \frac{1}{3} \pm z)$, and $(\frac{1}{3}, \frac{2}{3}, \frac{2}{3} \pm z)$. A straightforward interpretation of the Patterson function led to the z coordinates of the bromine atom and the carbon atom bonded to it. These two parameters and the isotropic temperature factor B for each atom, of the form $\exp(-B \sin^2\theta \cdot \lambda^{-2})$, were refined by least squares. The function minimized in the least squares calculation was $\sum w(|F_o| - |F_c|)^2 / \sum w|F_o|^2$, where w is the weighting factor and F_o and F_c are the observed and calculated structure factors. Atomic scattering factors for neutral bromine and carbon were taken from *International Tables for X-ray Crystallography* (Ibers, 1962). The

values for bromine were corrected for dispersion by subtracting 1.0 electron (Templeton, 1962). The imaginary part of the dispersion correction was neglected.

After four cycles the residual R value, defined as $R = \sum ||F_o| - |F_c|| / \sum |F_o|$, was 0.443. A three-dimensional electron density map was then calculated with use of signs determined by the carbon and bromine atoms. The resulting map permitted the evaluation of the parameters of the remaining seven carbon atoms in the asymmetric unit. The parameters of these nine atoms were refined by four cycles of least squares. This reduced R to 0.161. Bromine was then assigned an anisotropic temperature factor; four more cycles of least-squares refinement reduced R to 0.156.

A detailed comparison was then made between calculated and observed $|F|$ values. Four reflections had large discrepancies, and many more had discrepancies which seemed outside experimental error. As a result the intensities of 125 reflections were remeasured on the same crystal. Outside of the four obvious errors, no significant discrepancies between measured intensities were found. After correction of the intensities of the four reflections the bromine and eight carbon atoms were given anisotropic temperature factors and their parameters refined with four cycles of least squares. This reduced R to 0.139. The temperature factors for all atoms were satisfactorily small except for the B_{11} and B_{22} terms of bromine, which were 5.3, and the B_{33} term of C(4), which was 6.65.

The positions of the five hydrogen atoms were then calculated and included in the next refinement. The hydrogen coordinates were held fixed but the isotropic temperature factors were allowed to vary. This reduced R to 0.128.

Since R still seemed high, three least-squares refinements were calculated in which a portion of the experimental data was eliminated. In the first calculation zero weight was given to those reflections which had an appreciable discrepancy between calculated and observed structure factors. Thirty-two reflections were selected, and after two cycles of least-squares refinement R was 0.113. In the second and third calculations all low angle data below $\sin \theta / \lambda = 0.30$ and all high angle data above $\sin \theta / \lambda = 0.531$ were eliminated, respectively. After two cycles of least-squares refinement R became 0.116 when the low angle data were not used and 0.124 when the high angle data were not used. Although the first two procedures led to a small improvement in R the improvement did not seem sufficiently dramatic to warrant eliminating a portion of the experimental data.

Fourier calculations of the electron density showed a peak at $0, 0, 0.46$ about half as high as a carbon atom. Because of space limitations this position is impossible for any plausible chemical impurity which we have been able to invent. Attempts to identify an error in the data which would produce the peak also failed. We suggest that this peak is the result of disorder of

some kind, but we have failed to find a completely satisfactory model of the details of the disorder.

Least-squares calculations with one-tenth of a bromine atom, or with a carbon atom at that position, improved the agreement slightly and at the same time shifted other atoms by small amounts. The differences in coordinates found in these refinements suggest that standard deviations of positions should be stated as about 0.02 Å, rather than the values given in Table 1, to include the uncertainty introduced by our ignorance of the nature of the disorder.

Since no significant improvement could be effected by the procedures described above and no appreciable discrepancies could be found in the experimental data it was decided to terminate the refinement.

The positional and thermal parameters and their calculated standard deviations are listed in Tables 1 and 2 respectively for all atoms except hydrogen. These parameters were obtained with use of all data, not refining the positional parameters of the hydrogen atoms, and holding the temperature parameters of the hydrogen atoms equal, but allowing them to be refined. The *R* value of this refinement was 0.123. The observed and calculated structure factors are shown in Table 3.

Table 1. *Final positional parameters and standard deviations* in parentheses × 10⁴*

	<i>x</i>	<i>y</i>	<i>z</i>
Br	0.0	0.0	0.1052 (1)
C(1)	0.0	0.0	0.2173 (7)
C(20)	0.0	0.0	0.3595 (8)
C(2)	0.1212 (7)	0.1184 (7)	0.2490 (5)
C(3)	0.2241 (8)	0.2153 (9)	0.2099 (6)
C(4)	0.3275 (9)	0.3147 (8)	0.2525 (7)
C(5)	0.3249 (8)	0.3146 (8)	0.3291 (6)
C(6)	0.2207 (8)	0.2154 (9)	0.3679 (6)
C(7)	0.1198 (7)	0.1170 (7)	0.3282 (5)
H(20)†	0.0	0.0	0.420
H(3)†	0.225	0.215	0.150
H(4)†	0.405	0.395	0.220
H(5)†	0.400	0.390	0.360
H(6)†	0.225	0.215	0.430

* The standard deviations listed are those estimated by least squares. Because of disorder, as described in the text, the uncertainty is judged to be about twice as great as indicated by these values.

† These parameters not refined.

Description of the structure

The structure consists of three non-polar molecular doublets arranged on the three threefold inversion axes

Table 2. *Anisotropic thermal parameters (Å²) and standard deviations in parentheses × 10²*

The temperature factors are in the form $\exp(-\sum \beta_{ij}h_ih_jb_ib_j/4)$, where b_ib_j are reciprocal cell lengths.

	<i>B</i> ₁₁	<i>B</i> ₂₂	<i>B</i> ₃₃	<i>B</i> ₁₂	<i>B</i> ₁₃	<i>B</i> ₂₃
Br	4.78 (7)	B11	3.73 (9)	B11/2	0.0	0.0
C(1)	1.72 (31)	B11	2.38 (53)	B11/2	0.0	0.0
C(20)	1.58 (30)	B11	2.94 (58)	B11/2	0.0	0.0
C(2)	1.31 (27)	1.46 (29)	3.78 (38)	0.74 (23)	-0.16 (25)	-0.12 (26)
C(3)	2.32 (35)	2.51 (3i)	4.48 (46)	1.33 (30)	-0.08 (32)	-0.53 (32)
C(4)	2.31 (36)	1.63 (32)	6.80 (61)	1.16 (29)	0.36 (37)	-0.39 (35)
C(5)	1.95 (34)	2.09 (34)	4.97 (47)	1.07 (29)	0.42 (32)	0.91 (32)
C(6)	1.62 (22)	2.97 (25)	4.51 (33)	1.38 (20)	0.17 (21)	0.88 (23)
C(7)	1.17 (27)	1.48 (29)	3.53 (36)	0.46 (23)	-0.15 (25)	0.18 (25)

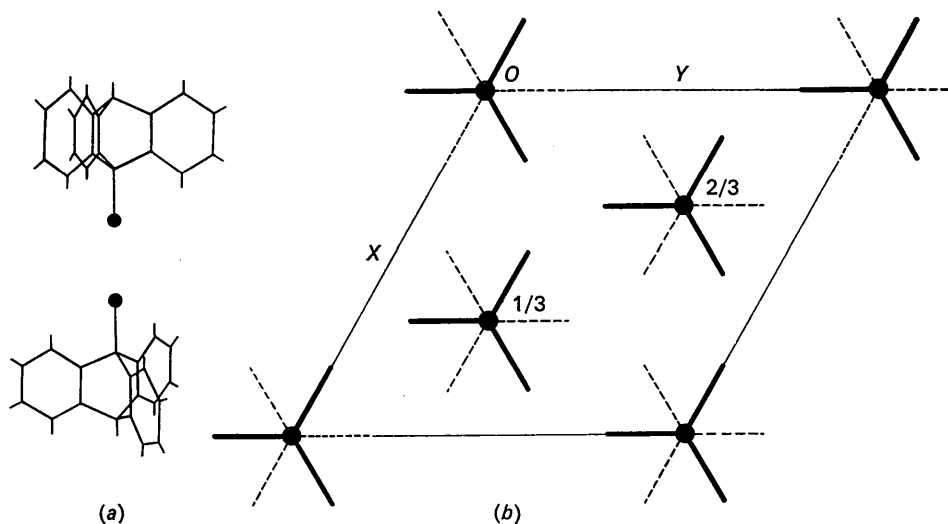


Fig. 3. (a) The stacking arrangement of two 1-bromotriptycene molecules about a center of symmetry. (b) The packing arrangement as viewed down the *c* axes. The numbers refer to the *z* coordinate of the inversion center.

Table 4. Interatomic distances with standard deviations ($\times 10^2$) for 1-bromotriptycene

Br—C(1)	1.97 (1) Å
C(1)—C(2)	1.53 (1)
C(2)—C(3)	1.37 (2)
C(2)—C(7)	1.39 (2)
C(3)—C(4)	1.42 (2)
C(4)—C(5)	1.35 (2)
C(5)—C(6)	1.39 (2)
C(6)—C(7)	1.37 (2)
C(7)—C(20)	1.51 (2)

Table 5. Interatomic bond angles with standard deviations ($\times 10$) for 1-bromotriptycene

Br—C(1)—C(2)	111.4 (6)°
H(20)—C(20)—C(7)	111.4 (6)
C(1)—C(2)—C(3)	128.5 (9)
C(20)—C(7)—C(6)	128.0 (9)
C(2)—C(1)—C(2)	107.5 (6)
C(7)—C(20)—C(7)	107.4 (6)
C(1)—C(2)—C(7)	110.7 (8)
C(20)—C(7)—C(2)	112.2 (8)
C(2)—C(3)—C(4)	118.1 (9)
C(7)—C(6)—C(5)	120.0 (9)
C(3)—C(4)—C(5)	121.1 (9)
C(2)—C(7)—C(6)	119.9 (7)
C(4)—C(5)—C(6)	120.2 (8)
C(3)—C(2)—C(7)	120.7 (7)
C(2)—C(3)—H(3)	120.2 (8)
C(3)—C(4)—H(4)	116.7 (10)
C(4)—C(5)—H(5)	120.6 (8)
C(5)—C(6)—H(6)	118.1 (8)

to an aliphatic carbon atom. The C(1)—C(2) and C(7)—C(20) distances of 1.53 Å and 1.51 Å, respectively, are in good agreement with the expected value of 1.515 Å for an sp^2 — sp^3 carbon bond (Dewar & Schmeising, 1960). It is concluded, therefore, that there is no interaction between benzene rings in 1-bromotriptycene. Since the bond angles and interatomic distances are the same for both bridges (within the accuracy of this investigation), it appears that the bromine atom has a negligible effect on adjacent bond angles and bond distances.

Estlin & Karle (1967) recently reported the results of a crystal structure investigation of β -chloroethyltriptycene. Their bond angles and distances are in good agreement with the results reported in this investigation.

We wish to thank Professor A. Streitwieser for suggesting this problem, Dr E. Ziegler for checking the purity, Dr F. T. Jones for the determination of the optical properties, Mrs Kay S. Lee for taking the Weissenberg photographs, Professor A. Pabst for the pre-

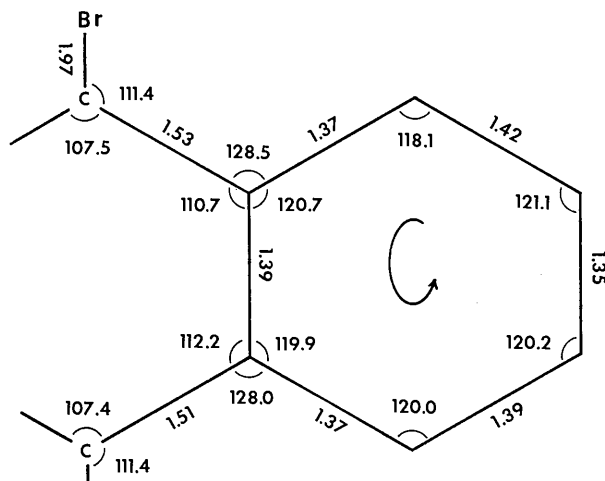


Fig. 4. Interatomic bond distances and angles for the asymmetric unit.

cession photographs and Dr A. Zalkin for many helpful suggestions and the use of his computer programs.

This work was supported in part by the U.S. Atomic Energy Commission.

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